

Lecture 12: RF Pulses - Basics

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Lectures 12 and 13 RF Pulses

- Lecture 12 (Oct. 3, 05)
 - Review the basics
 - A closer look of Bloch equation
 - Non-linearity of RF pulses
 - Flip angle and B1-field
 - RF pulse shapes
 - Rectangular pulses
 - Gaussian pulses
 - SINC pulses
 - Pulse functions
 - Excitation
 - Refocusing
 - inversion
- Lecture 13 (Oct 5, 05)
 - SLR pulses
 - Spatial selection and saturation
 - Spectral selection and saturation
 - Spatial-spectral selection
 - Magnetization Transfer
 - Excitation
 - Inversion
 - Refocusing

Reference: Part II, MRI Handbook by Bernstein, King and Zhou,

Review

- What's an RF pulse?
- Bloch equation:

$$\left(\frac{d\vec{M}}{dt} \right) = \gamma \vec{M} \times [\vec{B}_1(t) + \hat{z}B_0]$$

- Rotating frame:

$$\left(\frac{d\vec{M}}{dt} \right)_{rot} = \gamma \vec{M} \times \left[B_1(t) (\hat{x} \cos(\omega_f - \omega_0)t - \hat{y} \sin(\omega_f - \omega_0)t) + \hat{z} \left(B_0 - \frac{\omega_0}{\gamma} \right) \right]$$

- Resonance:

$$\omega_0 = \omega_f = \omega_{Lamor} = \gamma B_0$$

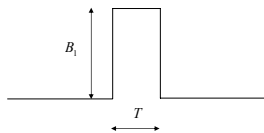
- Bloch equation on resonance:

$$\left(\frac{d\vec{M}}{dt} \right)_{rot} = \gamma \vec{M} \times [\hat{x}B_1(t)]$$

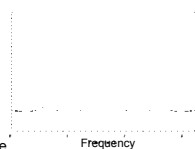
Let's think ...

- Can $\omega_f = \omega_0$ be always satisfied?
- Can $\omega_{Lamor} = \omega_0$ be always satisfied?

Consider a "Rectangular Pulse"



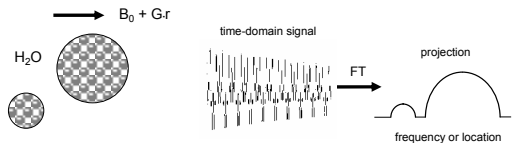
What are its frequency components?



Thus, $\omega_f = \omega_0$ cannot be always true.

RF (or B1) Off-Resonance

Consider a Spin System Subject to a Gradient



What is ω_{Lamor} ??

$$\omega_{Lamor}(x) = \gamma (B_0 + Gx)$$

Thus, $\omega_{Lamor} = \omega_0$ cannot be always true.

B0 Off-Resonance

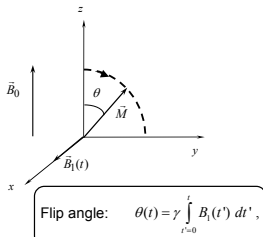
Non-linear Response of RF Pulses

- Off-resonance effect causes the Bloch equation to have non-linear responses.
- The frequency response of an RF pulse cannot be precisely obtained by FT; FT is only an approximation.

Bloch Equation in the RF Rotating Frame

$$\left(\frac{d\vec{M}}{dt}\right)_{rot} = \gamma \vec{M} \times \left[\hat{x}B_1(t) + \hat{z}\left(B_0 - \frac{\omega_{rf}}{\gamma}\right) \right]$$

Consider an RF Pulse



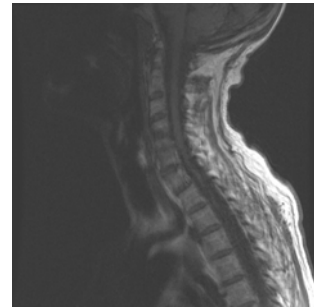
Example 3.1.1: A hard RF pulse has a rectangular shaped envelope. Its pulse width is 100µs, and its flip angle (on-resonance) is 90°. What is its amplitude for proton imaging?

Answer:

$$\theta = \frac{\pi}{2} = \gamma B_1 \times (100 \times 10^{-6} \text{ sec}),$$

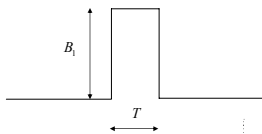
$$B_1 = \frac{\pi}{2\gamma(100\mu\text{s})} = \frac{\pi(\text{rad})}{2 \times 2\pi \times 42.57 \times 10^7 (\text{rad/sec} \cdot T) \times 100 \times 10^{-6} (\text{sec})} = 58.7 \mu\text{T}$$

B1-field inhomogeneity

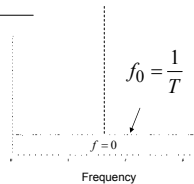


RF Pulse Shapes

- Rectangular (hard) pulse



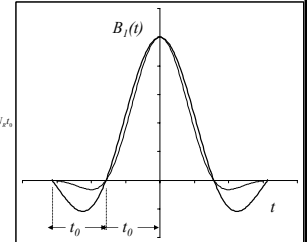
- Frequency Response
- Hard pulse is commonly used as non-selective pulse.
- RF un-blanking/blinking



SINC Pulses

- Truncated SINC and windowed (Hamming/Hanning) SINC

$$B_1(t) = \begin{cases} A t_0 \left[(1-\alpha) + \alpha \cos\left(\frac{\pi t}{N t_0}\right) \right] \frac{\sin\left(\frac{\pi t}{t_0}\right)}{\pi t}, & -N t_0 \leq t \leq N t_0 \\ 0, & \text{elsewhere} \end{cases}$$

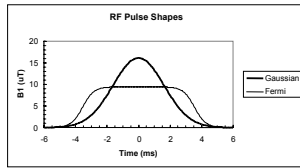


- Apodization
- Frequency response

$$Af = \frac{1}{t_0} = \frac{N_{lobe}}{T_{pulse_width}} \quad (\text{The central lobe is counted Twice})$$

- SINC pulse is commonly used as selective pulse.

Gaussian/Fermi Pulses



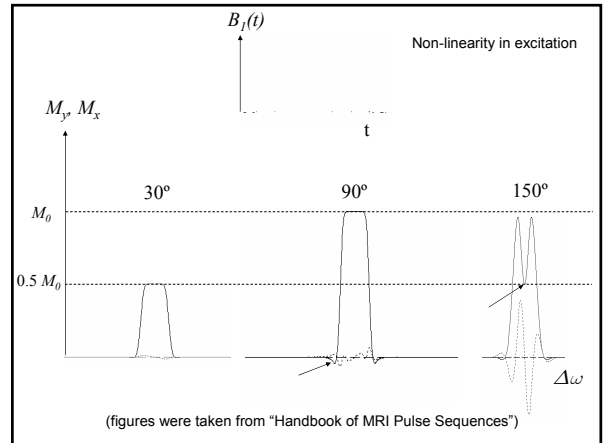
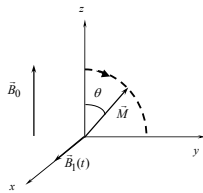
$$B_1(t) = A_G e^{-\frac{t^2}{2\sigma^2}} e^{i\Delta\omega_f t} \quad \text{Gaussian}$$

$$B_1(t) = \frac{A_F e^{i\Delta\omega_f t}}{1 + \exp\left(\frac{|t-a|}{a}\right)} \quad \text{Fermi}$$

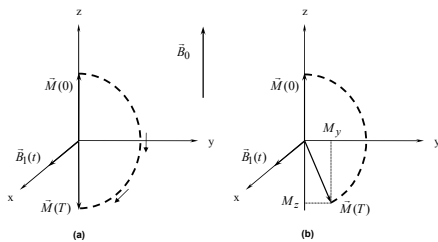
Tailored Pulse Shapes

- Shinnar-Le Roux Pulses
- Adiabatic Pulses
- ...

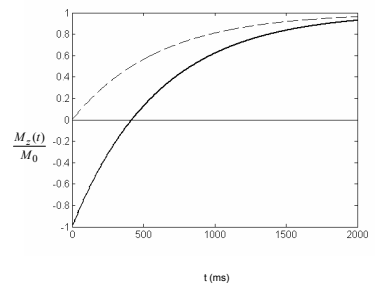
RF Pulse Functions - Excitation



Inversion



Saturation Recovery and Inversion Recovery



$$M_z(t) = M_0 \left[1 - (1 - \cos\theta_{inv}) e^{-t/T1} \right]$$

